

Augmented Lagrangian Preconditioner for Linear Stability Analysis of incompressible fluid flows on large configurations

J.Moulin', J-L. Pfister ${ }^{1}$, O.Marquet ${ }^{1}$, P. Jolivet²<br>I Office National d'Etudes et de Recherches Aérospatiales<br>${ }^{2}$ ENSEEIHT - Institut de Recherche en Informatique de Toulouse

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## Introduction

## What is Linear Stability Analysis ?

One wants to know if some steady solution of equation (1) is temporally stable or unstable:

$$
\begin{equation*}
\frac{\partial \boldsymbol{q}}{\partial t}=\boldsymbol{\mathcal { R }}(\boldsymbol{q}) \tag{1}
\end{equation*}
$$

## Method : Linear Stability Analysis

Step 1: compute a steady solution

$$
\mathcal{R}\left(\boldsymbol{q}_{\boldsymbol{b}}\right)=\mathbf{0}
$$

Step 2: test its stability for small monochromatic perturbations $\widehat{\boldsymbol{q}}(x) e^{\sigma t}$ around the steady solution $\boldsymbol{q}_{\boldsymbol{b}}$


## Introduction

## A classical example



From [Goharzadeh \& Molki, 2015]

## Typical question :

What is the critical Reynolds number above which the von Karman vortex street appears?

## Introduction

## A classical example

Step 1: compute a steady solution


Steady Navier-Stokes solution ( $R e=50$ ) [Sipp et al, 2010]

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$$
\begin{aligned}
\sigma M \widehat{\boldsymbol{q}} & =J\left(\boldsymbol{q}_{\boldsymbol{b}}\right) \widehat{\boldsymbol{q}} \\
\lambda & =\mathfrak{R}[\sigma] \\
\omega & =\Im[\sigma]
\end{aligned}
$$



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Steady Navier-Stokes
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[Sipp et al, 2010]

## Introduction

## Why Linear Stability Analysis?

Some nice features:

- Easy to determine a threshold value (sign of $\mathfrak{R}[\sigma]$ )
- Less expensive than nonlinear time-integration

But some computational burdens:

- Find a (not necessarily stable) steady solution: Newton method > Multiple inversions of $J$
- Find internal eigenvalues of generalized EV problems : Krylov-Schur + shift-and-invert
$>$ Multiple inversions of $\mathrm{J}-s M$, where $s$ is the shift


## Introduction

How to invert matrix of the type J - s M efficiently?
> For reasonably small configurations : direct sparse solvers (MUMPS, SUPERLU, etc)
$>$ For large configurations : iterative method (GMRES, BiCGSTAB, ...) + good preconditioner

## Introduction

Linearized incompressible Navier-Stokes operator (i.e. J -s M):


Once discretized with FE : classical saddle-point problem

$$
\left(\begin{array}{cc}
A & B^{T} \\
B & 0
\end{array}\right)\binom{u}{p}=\binom{f}{g}
$$

## How to precondition this?

- SIMPLE [Patankar 1980]
- Stokes Preconditioner [Tuckerman, 1989] (based on adaptation of existing time-stepping code)
- Pressure Convection Diffusion [Silvester et al. 2001]
- Least-Squares Commutator [Elman et al. 2006]
- Augmentated Lagrangian [Benzi and Olshanskii 2006], [Heister and Rapin 2013]


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## Overview

1 - Augmentation-based preconditioners

2 - Performances

3 - FreeFem++ parallel implementation

4 - Parallel 3D numerical examples

5 - Some further refinement ...

## 1-Augmentation-based preconditioners

## Augmented problems

## Augmented Lagrangian (algebraic augmentation)

$$
\left(\begin{array}{cc}
A & B^{T} \\
B & 0
\end{array}\right)\binom{u}{p}=\binom{f}{g} \quad \longrightarrow \quad\left(\begin{array}{cc}
A_{\gamma} & B^{T} \\
B & 0
\end{array}\right)\binom{u}{p}=\binom{f_{\gamma}}{g} \quad \begin{aligned}
& A_{\gamma}=A+\gamma B^{T} W^{-1} B \\
& f_{\gamma}=f+\gamma B^{T} W^{-1} g
\end{aligned}
$$

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\end{aligned}
$$

## Grad-Div augmentation (variational augmentation)

$$
\begin{gathered}
\int_{\Omega} s \boldsymbol{u} \cdot \breve{\boldsymbol{u}}+(\boldsymbol{u} \cdot \nabla) \boldsymbol{u} \cdot \check{\boldsymbol{u}}+R e^{-1} \nabla \boldsymbol{u}: \nabla \check{\boldsymbol{u}}-\mathrm{p} \boldsymbol{\nabla} \cdot \check{\boldsymbol{u}} \\
+\int_{\Omega} \gamma(\nabla \cdot \boldsymbol{u})(\nabla \cdot \check{\boldsymbol{u}})=\mathbf{0} \\
\quad-\int_{\Omega}(\nabla \cdot \boldsymbol{u}) \check{\boldsymbol{q}}=\mathbf{0} \\
\hline
\end{gathered}
$$

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Augmented Lagrangian leaves the discrete solution unchanged
Grad-Div leaves the continuous solution unchanged

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## 1- Augmentation-based preconditioners

Classical vs. modified version

In both cases, the same block structure arises :

$$
\left(\begin{array}{cc}
A_{\gamma} & B^{T} \\
B & 0
\end{array}\right)=(\underbrace{\left(\begin{array}{cc}
I & 0 \\
B A_{\gamma}^{-1} & I
\end{array}\right)} \underbrace{\left(\begin{array}{cc}
A_{\gamma} & 0 \\
0 & S
\end{array}\right)} \underbrace{\left(\begin{array}{cc}
I & A_{\gamma}^{-1} B^{T} \\
0 & I
\end{array}\right)} \quad S=-\mathrm{BA}_{\gamma}^{-1} \mathrm{~B}^{\mathrm{T}}
$$

Classical preconditioner

$$
\mathcal{P}_{\text {class }}=D U=\left(\begin{array}{cc}
A_{\gamma} & B^{T} \\
0 & S
\end{array}\right)
$$

with

$$
S^{-1} \simeq\left(\mathrm{Re}^{-1}+\gamma\right) M_{p}^{-1}-s\left(B M_{u} B^{T}\right)^{-1}
$$

$$
A_{\gamma}^{-1} \simeq \text { it's complicated ... }
$$

Main features:
> Mesh optimality
> Reynolds optimality
$>$ The higher $\gamma$, the less iterations ( $A_{\gamma}^{-1}$ ouch!)

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## Classical preconditioner

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\begin{gathered}
\mathcal{P}_{\text {class }}=D U=\left(\begin{array}{cc}
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\text { with }
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$S^{-1} \simeq\left(\mathrm{Re}^{-1}+\gamma\right) M_{p}^{-1}-s\left(B M_{u} B^{T}\right)^{-1}$
$A_{\gamma}^{-1} \simeq i t$ 's complicated ...
Main features:
> Mesh optimality
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## Modified preconditioner

$$
\mathcal{P}_{\text {modif }}=\left(\begin{array}{cc}
{\left[\begin{array}{cc}
A_{11, \gamma} & A_{12, r} \\
0 & A_{22, \gamma}
\end{array}\right]} & B^{T} \\
0 & S
\end{array}\right)
$$

$$
\begin{aligned}
& S^{-1} \simeq\left(\mathrm{Re}^{-1}+\gamma\right) M_{p}^{-1}-s\left(B M_{u} B^{T}\right)^{-1} \\
& A_{i i, \gamma}^{-1} \simeq \text { off-the-shelf algebraic multigrid }
\end{aligned}
$$

## Main features:

> Mesh optimality
> Reynolds dependent
> Exists an optimal and case dependent $\gamma$

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Classical vs. modified version

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Classical precondilioner

$$
\mathcal{P}_{\text {class }}=D U=\left(\begin{array}{cc}
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with
$S^{-1} \approx\left(\operatorname{Re}^{-1}+\gamma\right) M_{p}^{-1}-s\left(B M_{u} B^{T}\right)^{-1}$
$A_{\gamma}^{-1} \simeq i t+s$ complicated....

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```
Main features:
* Mesh optimality
> Reynolds optimality
* The higher }\gamma\mathrm{ , the less iterations ( }\mp@subsup{A}{,}{-1}\mathrm{ - ouch!)
```


## Main features:

> Mesh optimality
> Reynolds dependent

- Exists an optimal and case dependent $\gamma$


## 2- Performances

## Choice of $\gamma$

The choice of a good $\gamma$ is determinant for the preconditioning efficiency!

Bright side : since the preconditioner is independent of the mesh
> Optimal $\gamma$ can be found on a coarse mesh

Dark side: Optimal $\gamma$ is problem and $R e-$ dependent


Figure : Influence of $R e \in[10,120]$ on optimal $\gamma$ for modified Grad-Div preconditioner

## 2- Performances

CPU time in Newton method

| Mesh | Velocity <br> DOFs | Pressure <br> DOFs | Full MUMPS |  |  |  | Modified Grad-Div |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  | Facto <br> $(\mathbf{m} \boldsymbol{s})$ | Reso <br> $(\mathbf{m} \boldsymbol{s})$ | tot/ndof <br> $(\boldsymbol{\mu s})$ | Facto <br> $(\mathbf{m s})$ | Reso <br> $(\mathbf{m s})$ | tot/ndof <br> $(\boldsymbol{\mu} \boldsymbol{s})$ |  |
| $32 \times 32$ | 9900 |  | 140 | 0 | 20 | 30 | 50 | 14 |  |
| $64 \times 64$ | 39000 | 5000 | 810 | 10 | 27 | 320 | 250 | 20 |  |
| $96 \times 96$ | 88000 | 11000 | 2250 | 40 | 33 | 840 | 580 | 21 |  |
| $256 \times 256$ | 623400 | 78200 | 34480 | 290 | 62 | 8090 | 4780 | 25 |  |

Averaged timings for 1 Newton iteration (2D lid-driven cavity, $R e=100, \gamma=0,1$ )

## 2- Performances

CPU time in Newton method

| Mesh | Velocity <br> DOFs | Pressure <br> DOFs | Fall MUMPS <br> $(\mathbf{m s})$ |  |  | Reso <br> $(\mathbf{m s})$ | tot/ndof <br> $(\boldsymbol{\mu s})$ | Facto <br> $(\mathbf{m s})$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  | 140 | 0 | 20 | 30 | 50 | Reso <br> $(\mathbf{m s})$ |
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| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  | Reso <br> $(\mathbf{m s})$ | tot/ndof <br> $(\boldsymbol{\mu s})$ | Facto <br> $(\mathbf{m s})$ | Reso <br> $(\mathbf{m s})$ | tot/ndof <br> $(\boldsymbol{\mu s})$ |  |  |
| $8 \times 8 \times 8$ | 9900 |  | 3,2 | 0,01 | 263 | 0,6 | 0,26 | 112 |  |
| $16 \times 16 \times 16$ | 39000 | 5000 | 295 | 0,3 | 2675 | 21 | 2,8 | 274 |  |

Averaged timings for 1 Newton iteration (3D lid-driven cavity, $R e=100, \gamma=0,1$ )


## 2- Performances

CPU time for eigenvalue computation

| Mesh | Velocity <br> DOFs | Pressure <br> DOFs | Full MUMPS |  | Modified Grad-Div |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | Fact [s] | Eig [s] | Fact [s] | Eig [s] <br> (it. inner GMRES) |  |  |
| $32 \times 32$ | 9890 | 1269 | 0,27 | 0,36 | 0,05 | $9(29)$ |  |
| $64 \times 64$ | 39306 | 4978 | 1,7 | 1,3 | 0,45 | $34(30)$ |  |
| $256 \times 256$ | 623482 | 78192 | 85 | 36 | 15 | $841(30)$ |  |
| Timings for computing 10 ev with ARPACK (2D lid-driven cavity, $R e=100, \gamma=0,1)$ |  |  |  |  |  |  |  |

## 2- Performances

CPU time for eigenvalue computation

| Mesh | Velocity <br> DOFs | Pressure <br> DOFs | Full MUMPS |  | Modified Grad-Div |  |
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|  |  |  | Eig [s] | Fact [s] | Eig [s] <br> (it. inner GMRES) |  |
| $8 \times 8 \times 8$ | 14739 | 729 | 8 | 2 | 1,6 | $35(24)$ |
| $16 \times 16 \times 16$ | 107811 | 4913 | 753 | 31 | 57 | $353(23)$ |

Timings for computing 10 ev with ARPACK (3D lid-driven cavity, $R e=100, \gamma=0,1$ )

```
2- Performances
What to remember ?
```

Iterative strategy will be faster than the direct solver when : time facto >> time solving
> For Newton method : always the case because the jacobian is new at each iteration
> For eigenvalue computation : true only for large configurations (3D typically)

## 2- Performances

Krylov subspace recycling techniques and eigenvalue computation

Idea : In Krylov-Schur + shift-invert, one has to perform many $(J-s M)^{-1}$ with the same matrix ! > Why not use Krylov subspace recycling from one linear solve to the next ?

## 2- Performances

Krylov subspace recycling techniques and eigenvalue computation

Idea : In Krylov-Schur + shift-invert, one has to perform many $(J-s M)^{-1}$ with the same matrix ! > Why not use Krylov subspace recycling from one linear solve to the next?


Figure : Effect of recycling during eigenvalue computation. Test case : 2D circular cylinder at $R e=50$.
Preconditioner : Modified Grad-Div with $\gamma=1$
Eigenvalue solver : ARPACK with shift-invert

## 3- Parallel implementation in FreeFem++ <br> PETSc/SLEPc interface (P. Jolivet)

Ingredient 1 : handle the preconditioner's block structure
PETSC solution : use of PCFIELDSPLIT preconditioner

## FreeFem++ interface :

```
fespace Wh(th,[P2,P2,P2,P1]); // full space
Wh [u,v,w,p];
Wh [b,bv,bw,bp] = [1.0, 2.0, 3.0, 4.0];
string[int] names(4);
names[0] = "xvelocity" ;
names[1] = "yvelocity" ;
names[2] = "zvelocity" ;
names[3] = "pressure" ;
// Set PETSc solver
set(A, sparams = " -ksp_type fgmres -pc_type fieldsplit -pc_fieldsplit_type multiplicative"
    + " -prefix_push fieldsplit_xvelocity_
    + " -ksp_type preonly -pc_type lu -pc_factor_mat_solver_package mumps"
    + " -prefix_pop"
    + " -prefix_push fieldsplit_yvelocity_" , ... ... ..
    fields = b[], names = names);
```


## 3- Parallel implementation in FreeFem++ <br> PETSc/SLEPc interface (P. Jolivet)

Ingredient 2 : provide a specific Schur complement approximation
PETSC solution: PCFieldSplitGetSubKSP(pc, \&nfields, \&subksp) KSPSetOperators(subksp[nfields-1], Sapprox, Sapprox)

FreeFem++ interface :

```
fespace Wh(th,[P2,P2,P2,P1]); // full space
fespace Qh(th,P1); // pressure space
Wh [u,v,w,p];
Wh [b,bv,bw,bp] = [1.0, 2.0, 3.0, 4.0];
string[int] names(4);
names[0] = "xvelocity" ;
names[1] = "yvelocity" ;
names[2] = "zvelocity" ;
names[3] = "pressure" ;
Qh pind;
pind[] = 1:pind[].n;
Wh [list, listv, listw, listp]= [0, 0, 0, pind]; // correspondance between Wh and Qh pressure DOFs
matrix[int] S(1);
S[0]=vSchur(Qh,Qh); // Schur complement approximation
// Set PETSc solver
set(A, sparams = "
    fields = b[], names = names, schurPreconditioner = S, schurList = list[]);
```


## 3- Parallel implementation in FreeFem++ <br> PETSc/SLEPc interface (P. Jolivet)

Ingredient 3: provide the inverse Schur complement approx. as a composition of two simple inverses

$$
S^{-1} \simeq\left(\mathrm{Re}^{-1}+\gamma\right) M_{p}^{-1}-s L_{p}^{-1}
$$

PETSC solution : use of PCCOMPOSITE preconditioner
FreeFem++ interface :

```
matrix[int] S(2);
S[0]=vMp(Qh,Qh); // pressure mass matrix
S[1]=vLp(Qh,Qh); // pressure laplacian matrix
// Set PETSc solver
set(A, sparams = " ... ... ... "
    + " -prefix_push fieldsplit_pressure_
        -ksp_type preonly -pc_type composite -pc_composite_type additive"
                            + " -prefix_push sub_0_"
                                + " -pc_type ksp -ksp_ksp_type cg -ksp_pc_type jacobi"
    + " -prefix_pop"
    + " -prefix_push sub_1_"
        + " -pc_type ksp -ksp_ksp_type fgmres -ksp_pc_type gamg"
        + " -prefix_pop"
    + " -prefix_pop" ,
    fields = b[], names = names, schurPreconditioner = S, schurList = list[]);
```


## 3- Parallel implementation in FreeFem++ <br> PETSc/SLEPc interface (P. Jolivet)

Ingredient 4: Recycling of Krylov basis bewteen two consecutive solve $(J-s M)^{-1}$ in SLEPC

PETSc solution : interface HPDDM's solvers with PETSc/SLEPc
FreeFem++ interface :

```
// Set SLEPc eigensolver Ax = sigma Bx
int k = zeigensolver
(DistA,
DistB,
vectors = EigenVEC, // Array to store the FEM-EigenFunctions
values = EigenVAL, // Array to store the EigenValues
sparams = " -eps_type krylovschur -st_type sinvert -eps_target 0+0.6i"
    + " -st_ksp_type hpddm -hpddm_st_krylov_method gcrodr -hpddm_st_variant flexible ... ",
    fields = b[], names = names, schurPreconditioner = S, schurList = list[]);
```


## 4- Parallel 3D numerical examples

Flow around low aspect-ratio flat plates [Marquet \& Larsson 2015]

Test case :
$>1$ million tetrahedrons / Taylor-Hood FE pair / 4,8 millions DOFs
$>R e=100$
Solvers:

- Steady solution : Newton method with FGMRES preconditioned by Modified Grad-Div ( $\gamma_{\text {optimal }}=0,3$ )
> velocity sub-blocks solved with FGMRES preconditionned by ASM, overlap=1, tol $=10^{-1}$
> Schur complement sub-block solved with CG preconditionned by jacobi, tol $=10^{-3}$
- Eigenvalues : Krylov-Schur + shift-invert + GCRO-DR(100,30) preconditioned by Modified Grad-Div ( $\gamma_{\text {optimal }}=$ $0,3)$
$>$ velocity sub-blocks solved with FGMRES preconditionned by ASM, overlap $=1$, tol $=10^{-1}$
> Schur complement sub-block 1 solved with CG preconditionned by jacobi, tol=10-3
> Schur complement sub-block 2 solved with FMGRES preconditionned by gamg, tol=10-3


## 4- Parallel 3D numerical examples

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$>$ velocity sub-blocks solved with FGMRES preconditionned by ASM, overlap=1, tol=10-1
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> Schur complement sub-block 2 solved with FMGRES preconditionned by gamg, tol=10-3


[^0] .


## 4- Parallel 3D numerical examples

## Flow around low aspect-ratio flat plates [Marquet \& Larsson 2015]

Test case :
$>1$ million tetrahedrons / Taylor-Hood FE pair / 4,8 millions DOFs
$\Rightarrow R e=100$

## Solvers:

- Steady solution : Newton method with FGMRES preconditioned by Modified Grad-Div ( $\gamma_{\text {optimal }}=0,3$ )
$>$ velocity sub-blocks solved with FGMRES preconditionned by ASM, overlap=1, tol=10-1
> Schur complement sub-block solved with CG preconditionned by jacobi, tol=10-3
- Eigenvalues : Krylov-Schur + shift-invert + GCRO-DR(100,30) preconditioned by Modified Grad-Div ( $\gamma_{\text {optimal }}=$ 0,3)
$>$ velocity sub-blocks solved with FGMRES preconditionned by ASM, overlap $=1$, tol $=10^{-1}$
> Schur complement sub-block 1 solved with CG preconditionned by jacobi, tol=10-3
> Schur complement sub-block 2 solved with FMGRES preconditionned by gamg, tol=10-3



## 4- Parallel 3D numerical examples

Flow around low aspect-ratio flat plates [Marquet \& Larsson 2015]


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Flow around low aspect-ratio flat plates [Marquet \& Larsson 2015]



Eigenvalue computation
( 10 ev requested with tolerance $10^{-6}$ )
The loss of scaling for high number of procs is mainly due to the non-optimality of ASM w.r.t. number of domains. To be improved ...

## 5- Some further refinement ...

Influence of $\gamma$ on the solution ?

## Grad-Div augmentation (variational augmentation)

$$
\begin{gathered}
\int_{\Omega} s \boldsymbol{u} \cdot \check{\boldsymbol{u}}+(\boldsymbol{u} \cdot \nabla) \boldsymbol{u} \cdot \check{\boldsymbol{u}}+R e^{-1} \nabla \boldsymbol{u}: \nabla \check{\boldsymbol{u}}-\mathrm{p} \nabla \cdot \check{\boldsymbol{u}} \\
+\int_{\Omega} \gamma(\nabla \cdot \boldsymbol{u})(\nabla \cdot \breve{\boldsymbol{u}})=\mathbf{0} \\
-\int_{\Omega}(\nabla \cdot \boldsymbol{u}) \breve{\boldsymbol{q}}=\mathbf{0}
\end{gathered}
$$

Grad-Div leaves the continuous solution unchanged But ... changes the discrete solution !


Figure : Eigenvalue spectrum of the flow around a 2D circular cylinder at $R e=50$

Spatial discretization: Taylor-Hood $\left(P_{2}, P_{1}\right)$

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Not a divergence free element !!
$\nabla \cdot\left(P_{2}\right) \notin P_{1}$

## 5-Some further refinement ...

## Influence of $\gamma$ on the solution ?

What if one uses a divergence-free element ?
$>$ Scott-Vogelius FE pair: $\left(P_{2}, P_{1}^{d c}\right)$ s.t. $\nabla \cdot\left(P_{2}\right) \in P_{1}^{d c}$


Figure : Eigenvalue spectrum of the flow around a 2D circular cylinder at $R e=50$

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Figure : Eigenvalue spectrum of the flow around a 2D circular cylinder at $R e=50$

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A few remarks :

- $\quad\left(P_{2}, P_{1}^{d c}\right)$ is inf-sup stable only on specific types of mesh (Hsieh-Clough-Toucher triangulation)
- We showed that when using divergence-free elements the variational and discrete augmentations are equivalent
- It is unprcatical to use the discrete augmentation without divergence free elements due to the unsparse nature of the augmentation term ...


## Conclusion

## Conclusion :

- Krylov subspaces iterative method preconditioned by Modified Grad-Div where shown to be efficient both for finding a steady solution and computing its spectrum
- Large 3D configurations and large number of processors accentuate the benefits of using the iterative strategy w.r.t. direct solver.
- Ritz vector recycling was shown to provide significant acceleration of the eigenvalue computation when using an iterative strategy for $(J-s M)^{-1}$
- A parallel implementation in FreeFem++/PETSc/SLEPc was proposed.


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## Perspectives:

- Scalings must be improved : find an optimal preconditioner for velocity sub-blocks
- Extension for preconditioning turbulence models (RANS equations)
- Towards coupled fluid-structure Linear Stability Analysis on large 3D configurations ...



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\text { Fluid-structure Jacobian matrix }=\left(\begin{array}{ll}
J_{f f} & J_{f s} \\
J_{s f} & J_{s s}
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$$
?
$$

## 2- Performances

Memory requirements in Newton method

| Mesh | Velocity <br> DOFs | Pressure <br> DOFs | Memory <br> direct MUMPS <br> $($ Mb) | Memory <br> Modified Grad-Div (Mb) | Memory gain (\%) |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $16 \times 16$ | 2600 | 340 | 5 | $2 \times 2$ | 20 |
| $32 \times 32$ | 9900 | 1300 | 16 | $2 \times 4$ | 50 |
| $64 \times 64$ | 39000 | 5000 | 75 | $2 \times 17$ | 55 |
| $96 \times 96$ | 88000 | 11000 | 191 | $2 \times 41$ | 57 |
| $256 \times 256$ | 623400 | 78200 | 1862 | $2 \times 353$ | 62 |

Memory requirements (2D lid-driven cavity, $R e=100, \gamma=0,1$ )

| Mesh | Velocity <br> DOFs | Pressure <br> DOFs | Memory <br> direct MUMPS <br> (Mb) | Memory <br> Modified Grad-Div (Mb) | Memory gain (\%) |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $8 \times 8 \times 8$ | 14700 | 729 | 143 | $3 \times 21$ | 56 |
| $16 \times 16 \times 16$ | 107800 | 4900 | 2565 | $3 \times 260$ | 70 |

Memory requirements (3D lid-driven cavity, $R e=100, \gamma=0,1$ )


[^0]:    Steady solution (axial velocity)

